



Nucleation and anomalous properties of pure and salty water

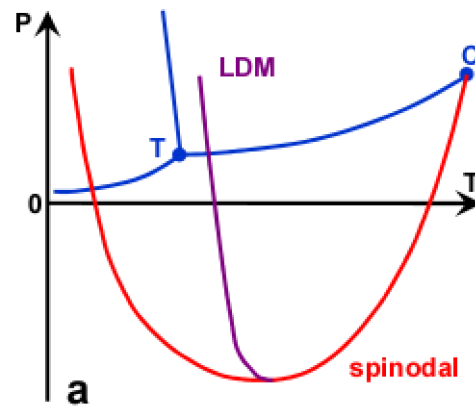
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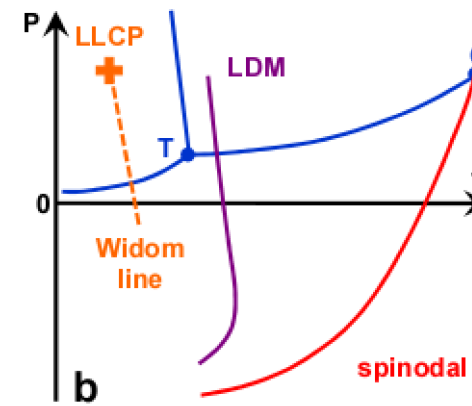
Motivation

In 1976 Speedy and Angell provided experimental evidence of a divergence in the isothermal compressibility of supercooled water

Possible Scenarios



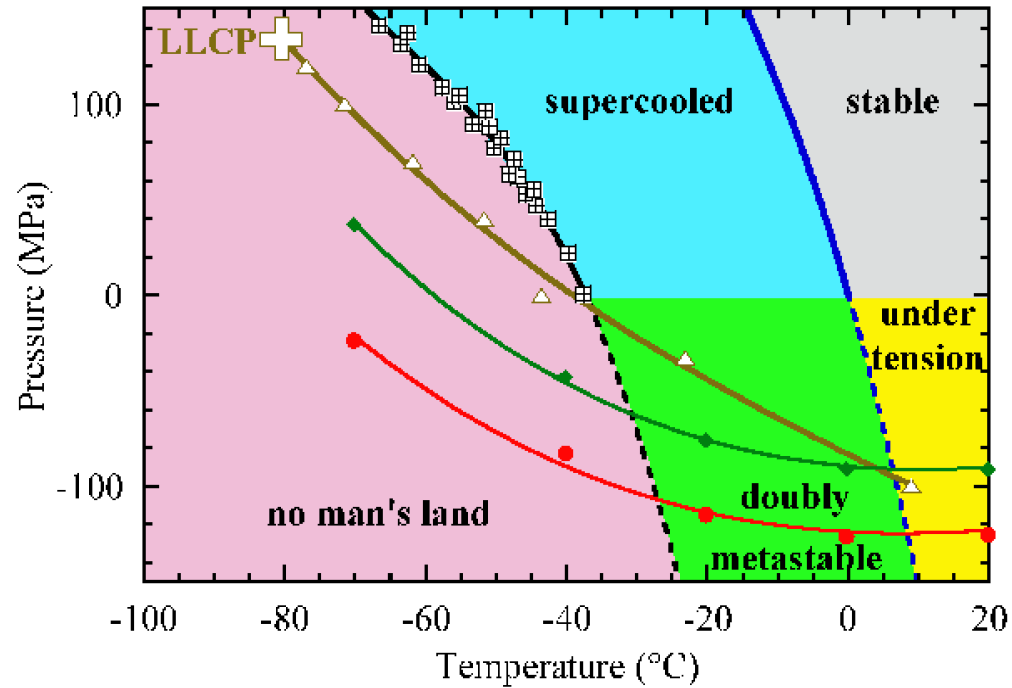
Stability limit conjecture
(retracing spinodal)



Liquid-liquid critical point
(retracing TMD)

The interesting region lies in the **no man's land**, not accessible to experiment because it is beyond the homogeneous nucleation line

The no man's land region

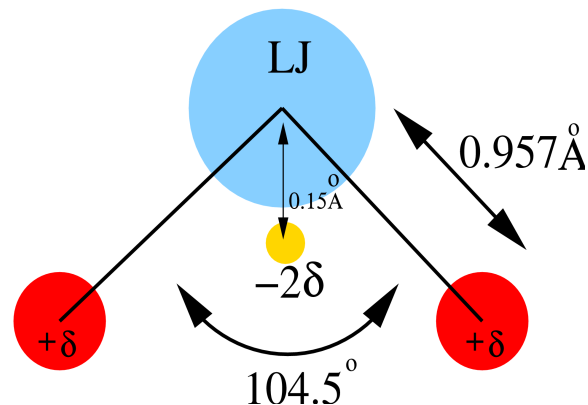


A difficult puzzle to solve

Because of the experimental difficulties (no man's land region),
numerical simulation may provide important information

but ...

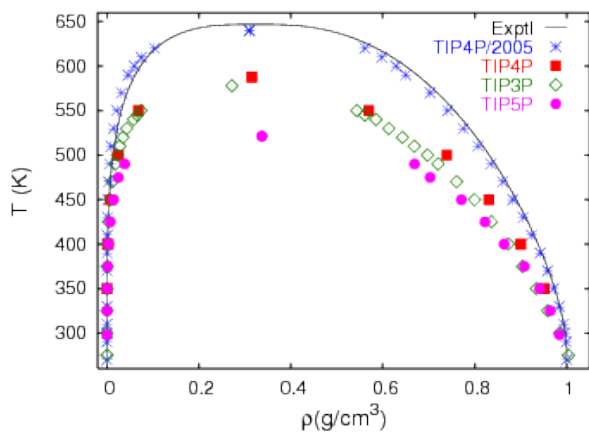
is the water force field accurate enough?



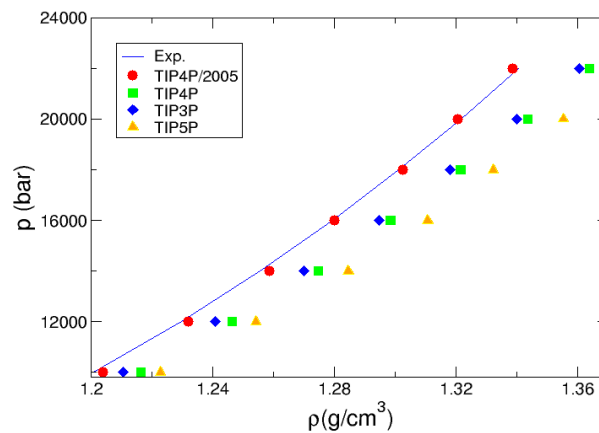
TIP4P/2005

Performance of TIP4P/2005

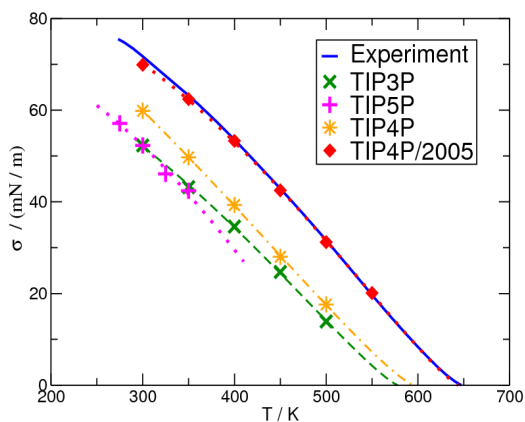
VLE



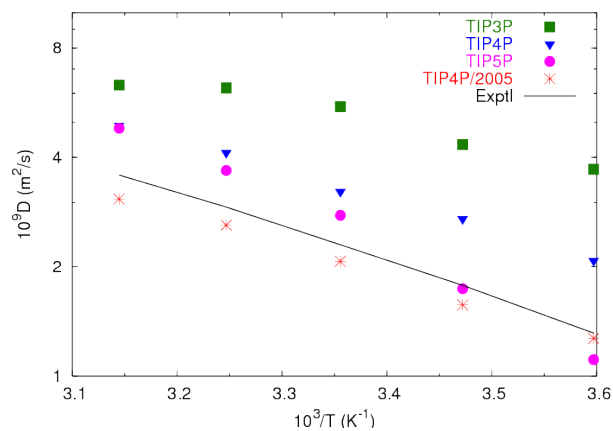
EOS high pressure



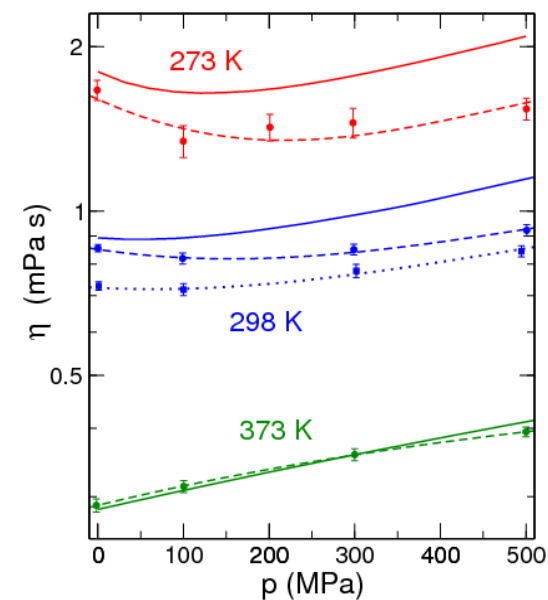
Surface Tension



Self Diffusion Coefficient



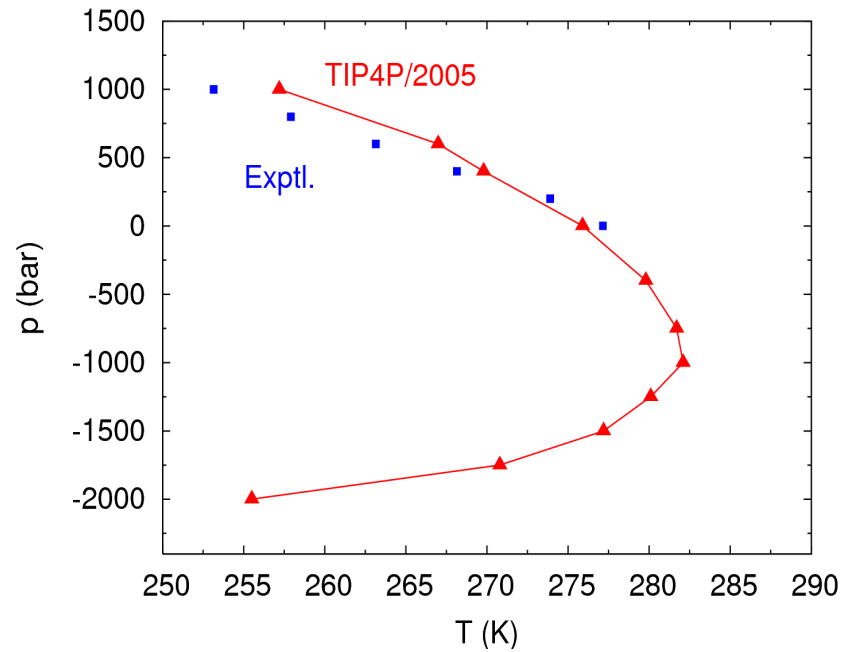
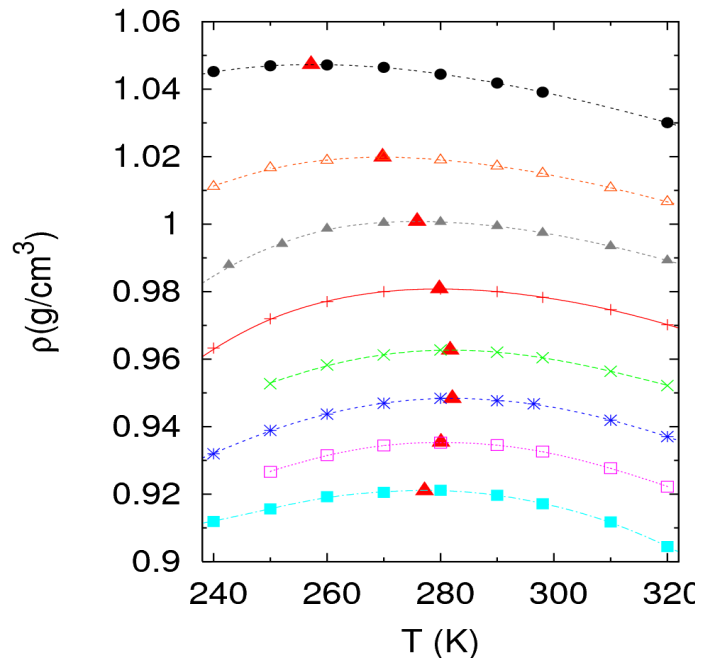
Shear Viscosity



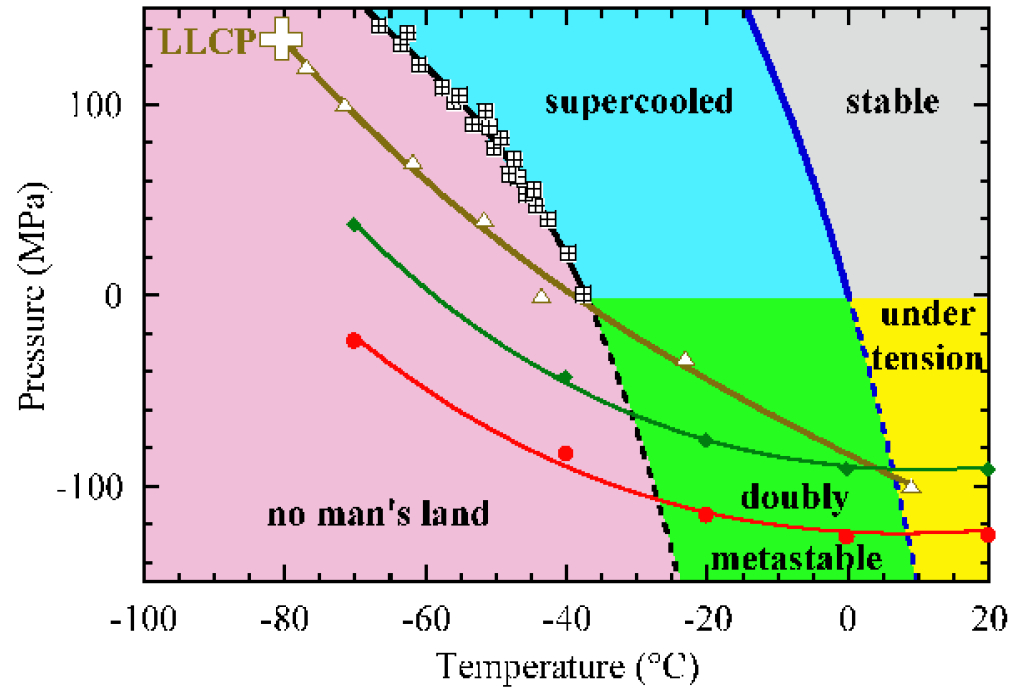
Property	TIP3P	SPC/E	TIP4P	TIP4P/2005	TIP5P
Enthalpy of phase change	4.0	2.5	7.5	5.0	8.0
Critical point properties	3.7	5.3	6.3	7.3	3.3
Surface tension	0.0	4.5	1.5	9.0	0.0
Melting properties	2.0	5.0	6.3	8.8	4.5
Liquid densities and TMD	1.8	5.5	4.0	8.5	4.0
Isothermal compressibil.	2.5	7.5	2.5	9.0	4.0
Gas properties	2.7	0.7	1.3	0.0	1.0
C_p	4.5	3.5	4.0	3.5	0.0
Dielectric constant	2.0	2.3	2.3	2.7	2.3
T_m/T_c , TMD- T_m	4.3	6.7	8.3	8.3	6.7
Densities of ices	3.5	5.0	6.0	8.8	2.3
EOS high pressure	7.5	8.0	7.5	10	5.5
Self diffusion coeff.	0.3	8.0	4.3	8.0	4.5
Shear viscosity	1.0	7.5	2.5	9.5	4.0
Orientational relaxation time	0.0	6.0	2.0	9.0	4.0
Structure	4.0	6.5	7.0	8.5	8.0
Phase diagram	2.0	2.0	8.0	8.0	2.0
Final score	2.7	5.1	4.7	7.2	3.7

Temperature of maximum density

(TIP4P/2005 vs experiment)



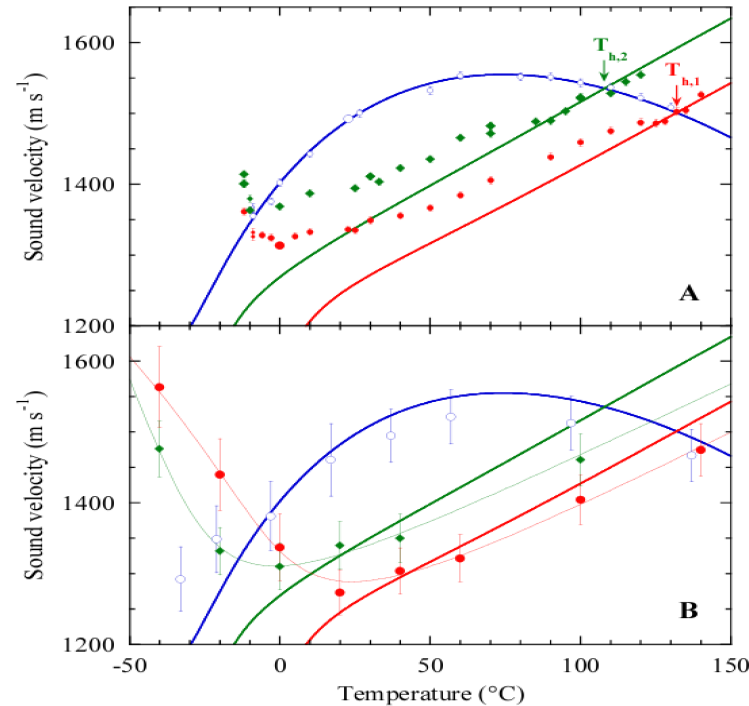
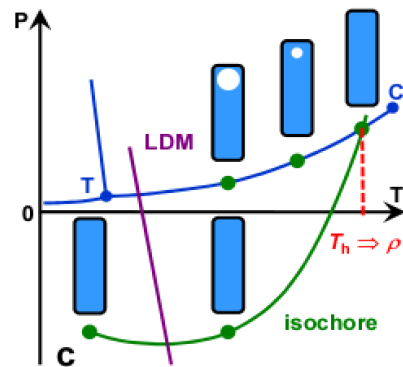
The doubly metastable region



It is possible to cross the (κ_T) Widom line before the system nucleates!

Speed of sound

$$c = \sqrt{\frac{1}{\rho \kappa_S}} = \sqrt{\frac{C_P}{C_V} \frac{1}{\rho \kappa_T}}$$



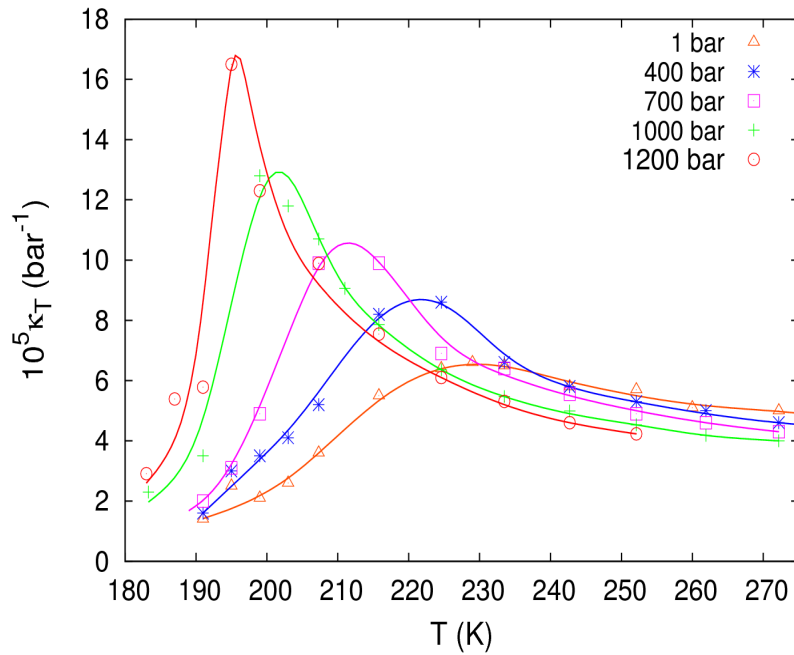
Experiment

Simulation

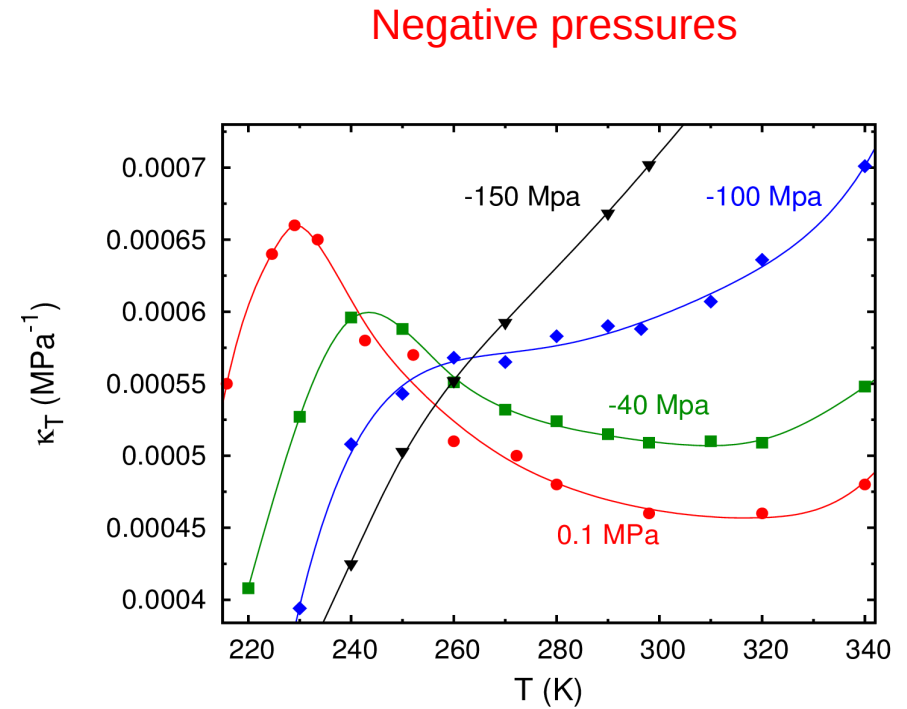
This is the first **experimental** report of extrema for a response function in supercooled **bulk** water!

Pallares et al., PNAS 111, 7936 (2014)

Isothermal compressibility

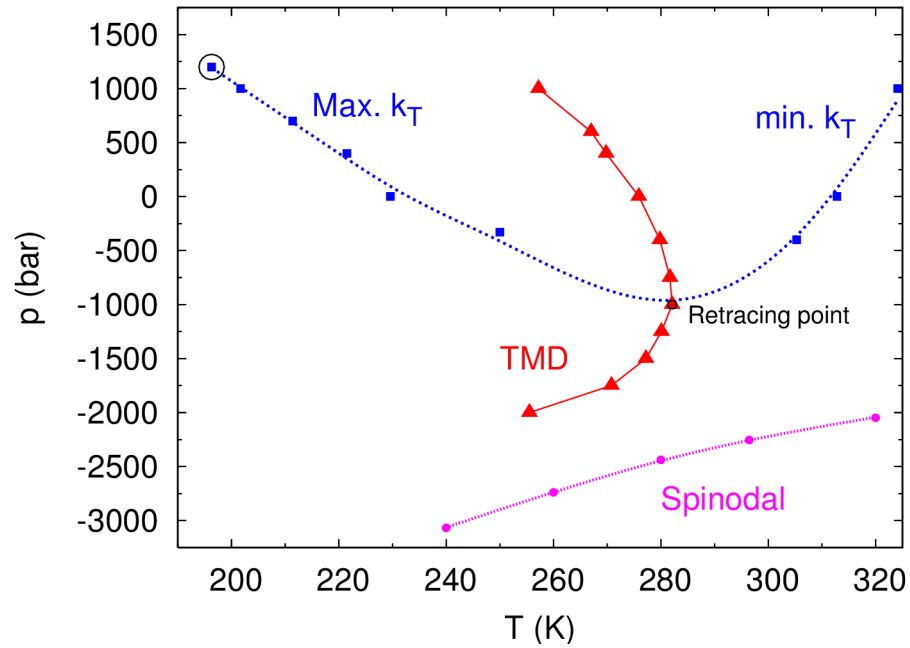


Positive pressures

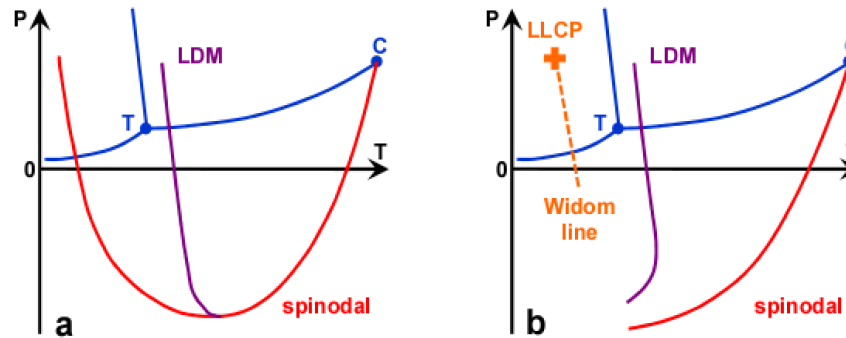
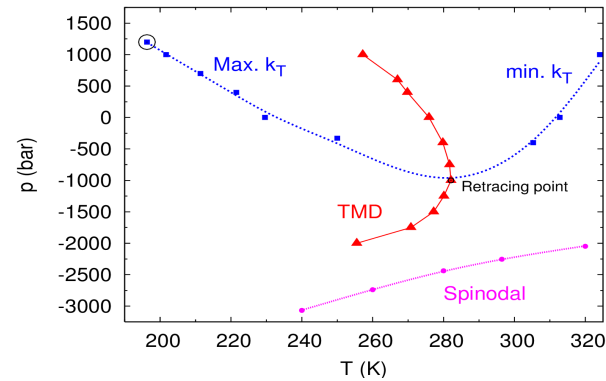


Maxima and minima collapse
below -100 MPa

Scenario of water anomalies for TIP4P/2005 (I)



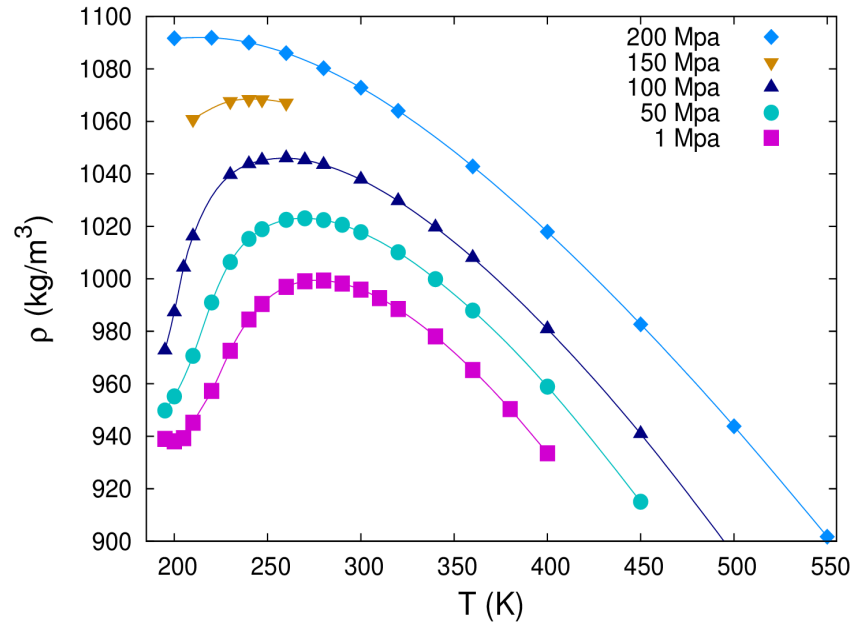
Scenario of water anomalies for TIP4P/2005 (I)



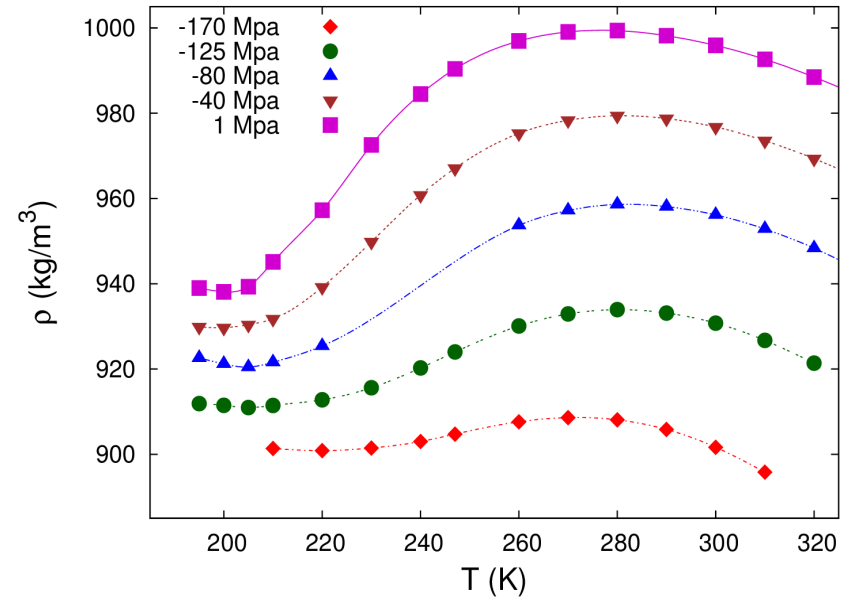
Singularity free scenario must be ruled out

TIP4P/2005 results compatible with the liquid-liquid critical point hypothesis

Temperature of minimum density (T_{mD})



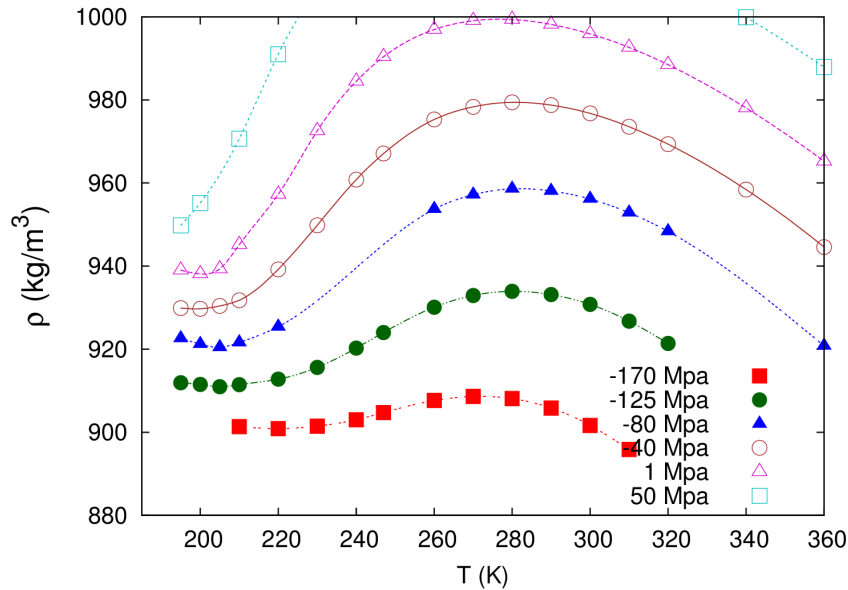
Positive pressures



Negative pressures

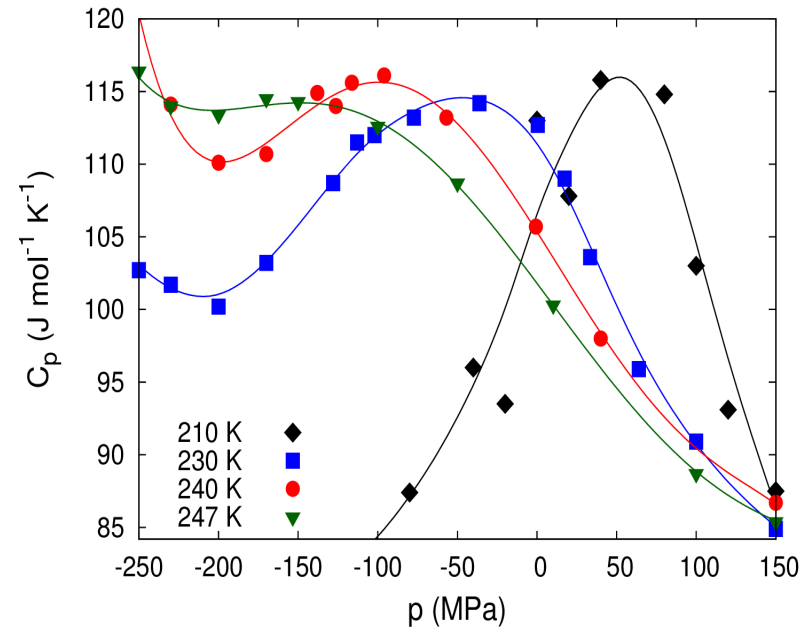
Density minima are featured in the stretched region

TmD and heat capacity extrema (C_p)

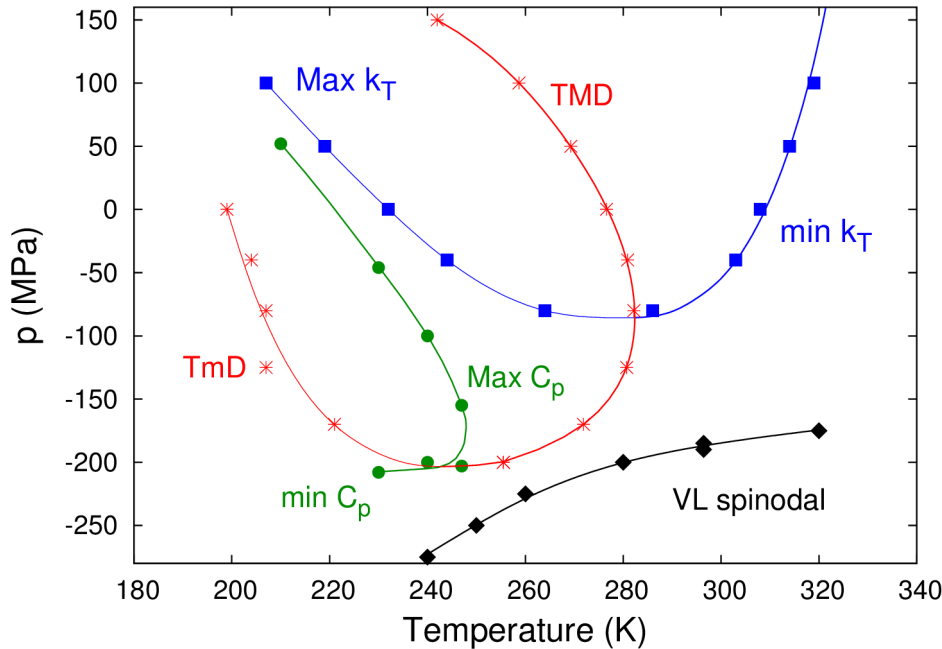
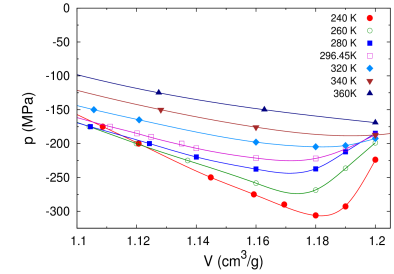
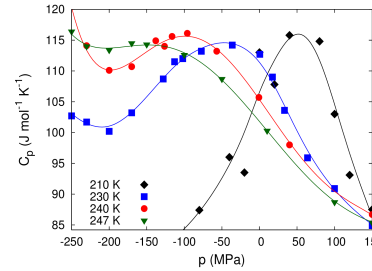
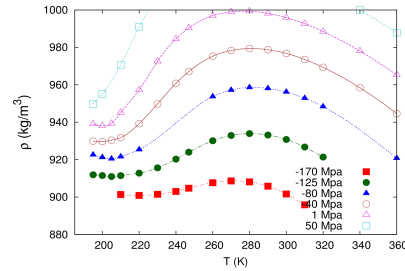
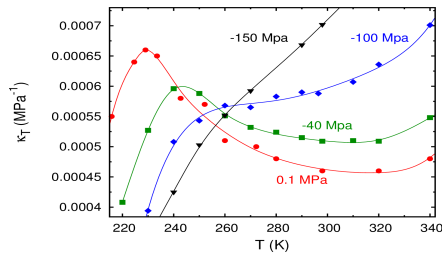


The maxima and minima of C_p appear only in the supercooled region (usually, in the doubly metastable region)

Density minima are featured in the stretched region

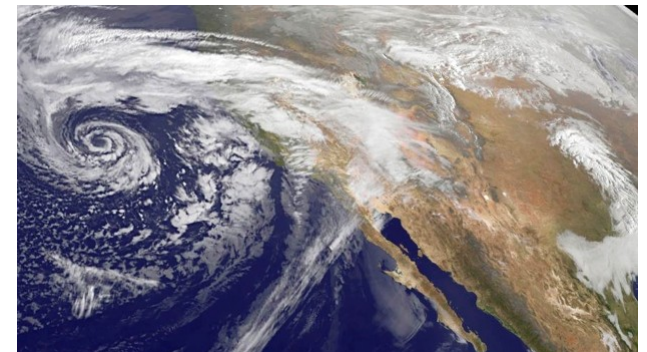
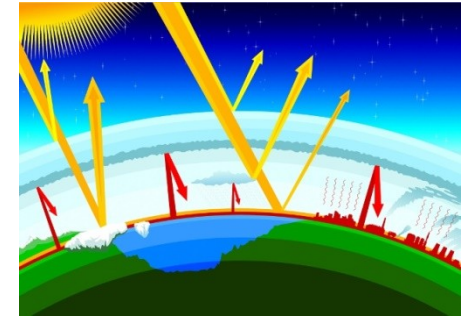
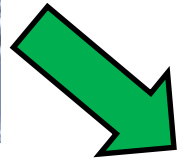
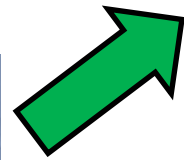
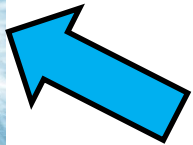


Full scenario of water anomalies (TIP4P/2005)

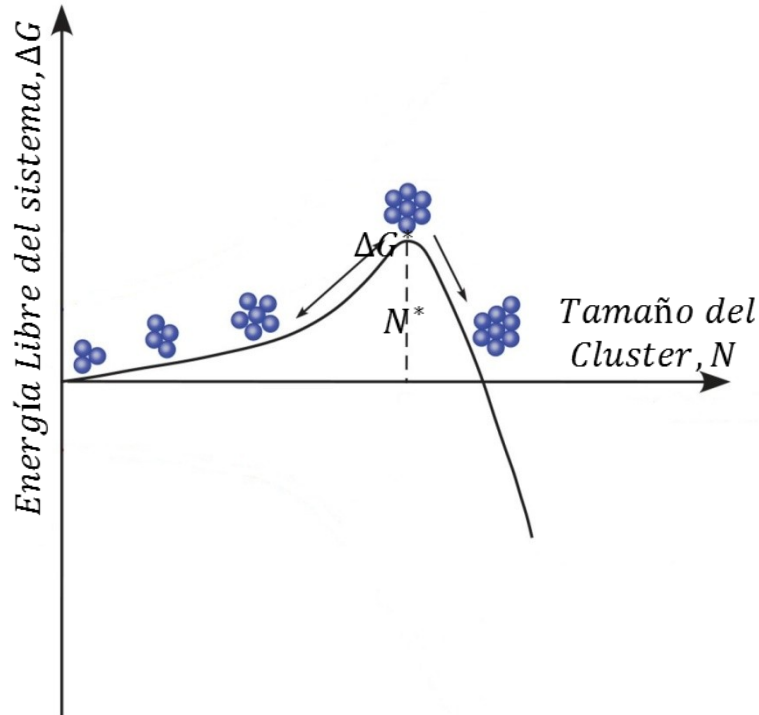


The liquid-liquid critical point scenario seems plausible but it could also be possible that $T_c = 0$ K

Relevance of ice nucleation



Homogeneous nucleation: free energy and nucleation rate



Classical nucleation theory (CNT):

$$\Delta G = -|\Delta\mu| \cdot N + A \cdot \gamma$$



$$\Delta G^* = \frac{|\Delta\mu| \cdot N^*}{2}$$

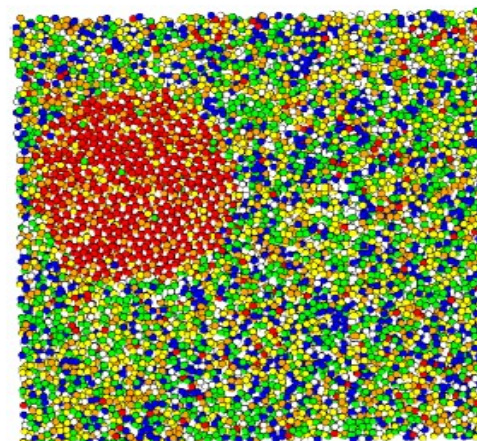
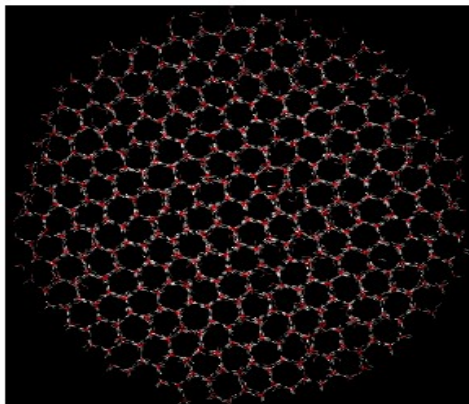
$$J = Z f^+ \rho_l \cdot e^{-\frac{\Delta G^*}{k_B T}}$$

Z: Zeldovich factor

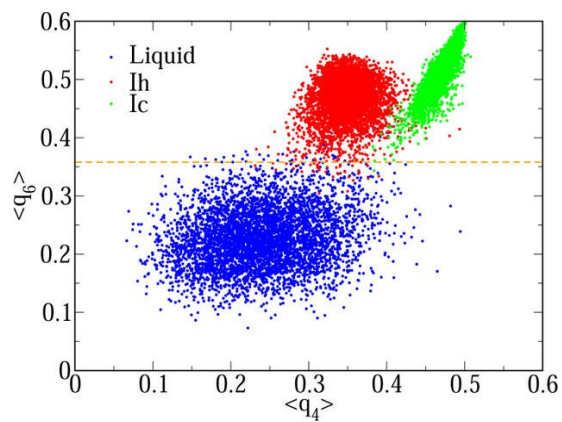
f⁺: attachment rate

ρ_l : liquid density

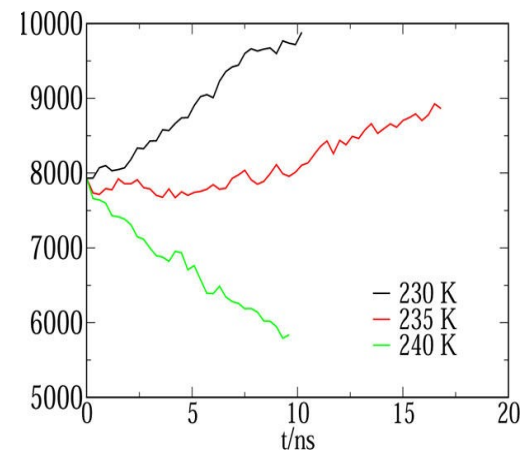
Seeding technique to evaluate N^*



Generation of a spherical ice cluster



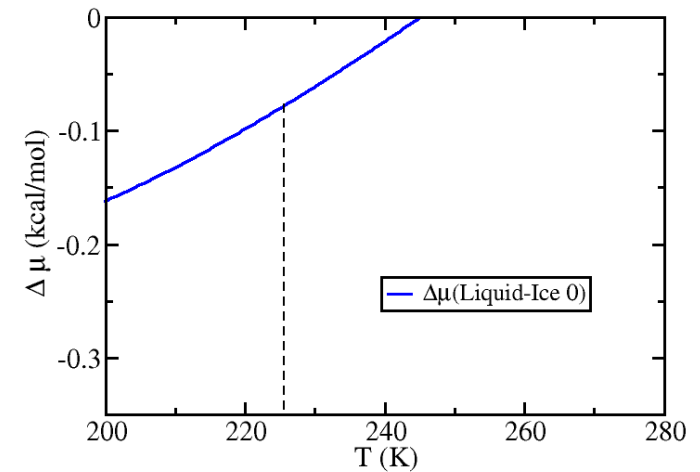
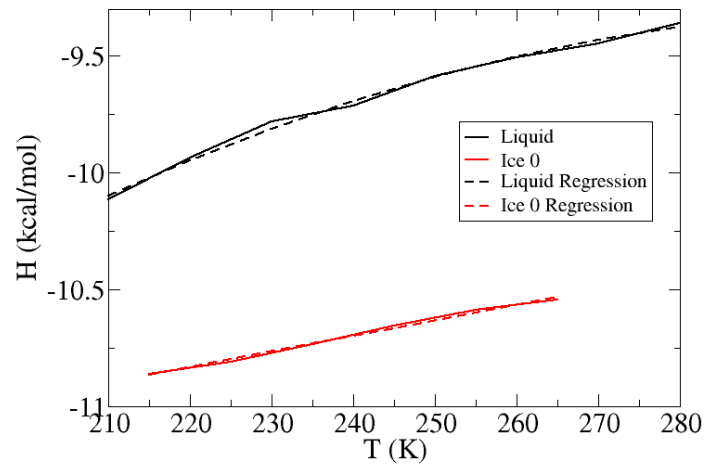
Insertion of the ice cluster in liquid water



Order parameter to detect the ice molecules

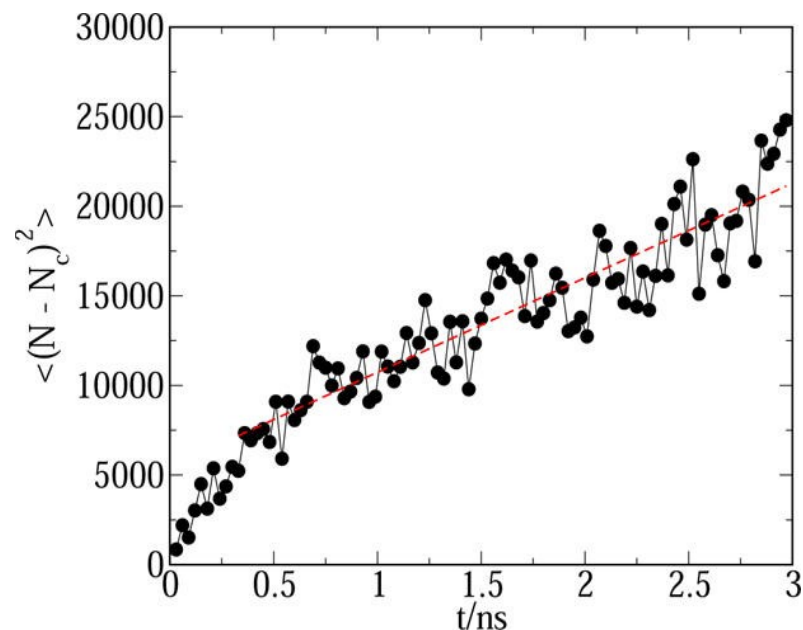
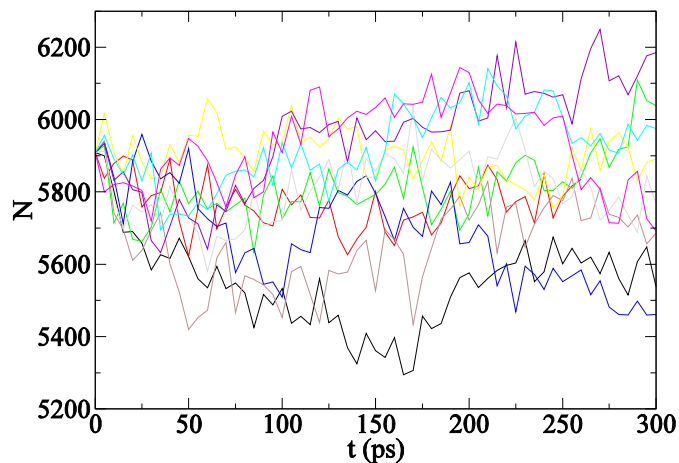
Molecular Dynamics simulations at different temperatures lead to bound the temperature at which the cluster is critical

Calculation of $\Delta\mu$: thermodynamic integration



$$\frac{\mu(T_2, p)}{k_B T_2} - \frac{\mu(T_1, p)}{k_B T_1} = - \int_{T_1}^{T_2} \frac{H(T)}{N k_B T^2} dT$$

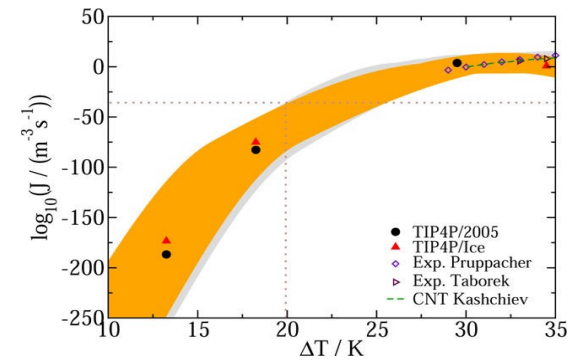
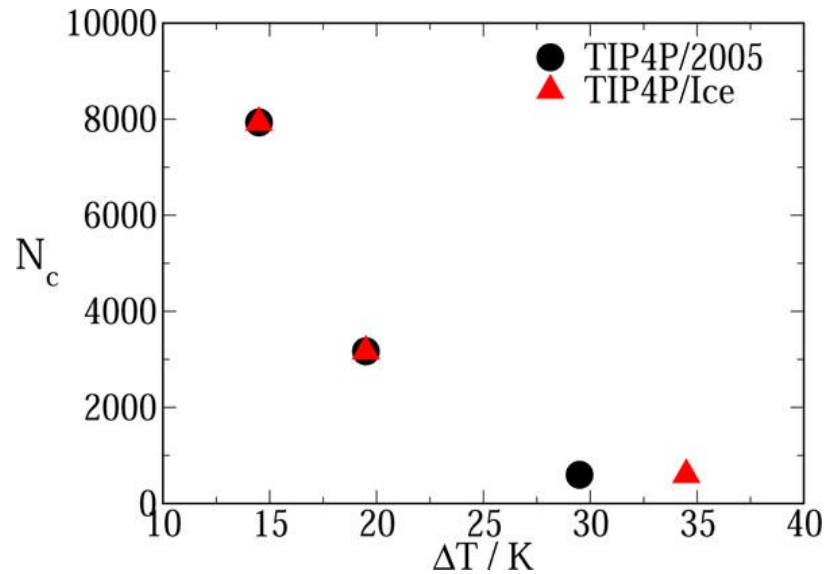
Calculation of the attachment rate f^+



Perform a number of simulations
of a cluster at the temperature
at which the cluster is critical

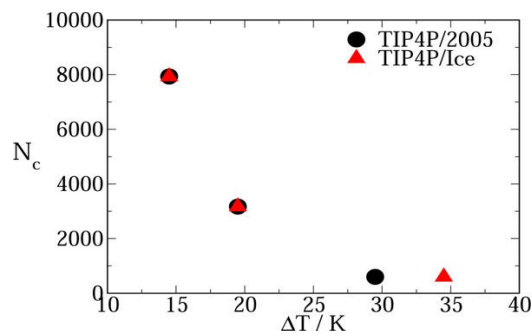
$$f^+ = \text{slope}/2$$

Relevant results: Size of the critical cluster and nucleation rates

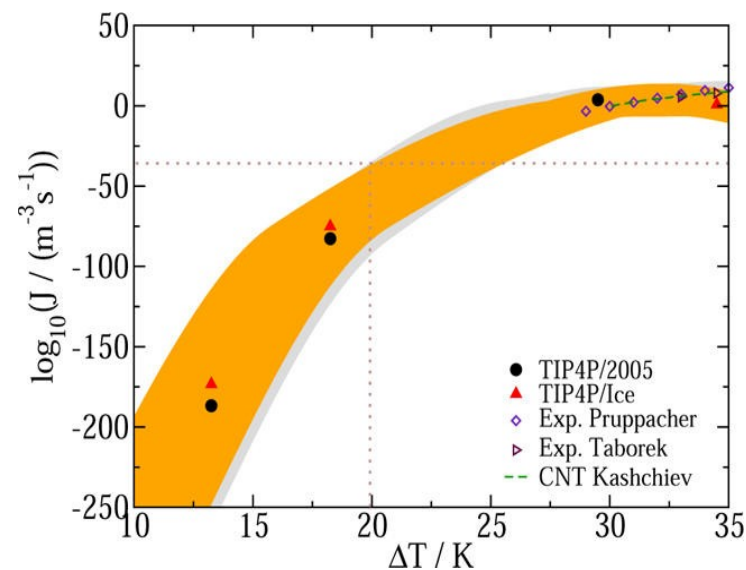


The size of the critical cluster is almost independent of the water model

Size of the critical cluster and nucleation rates



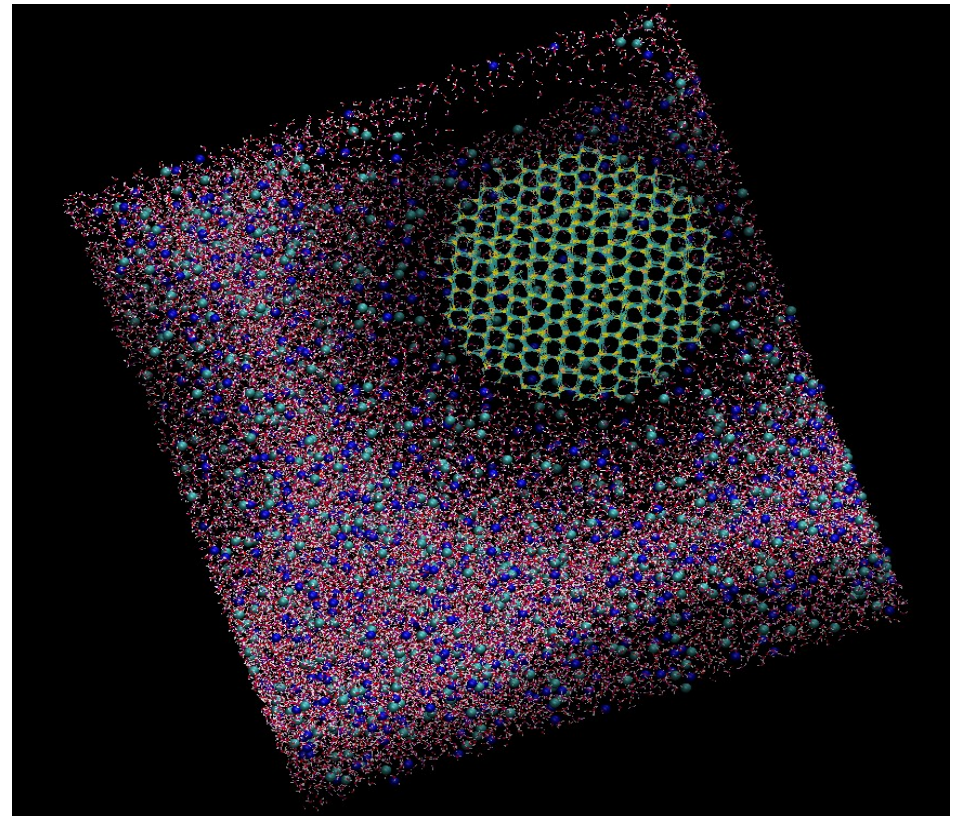
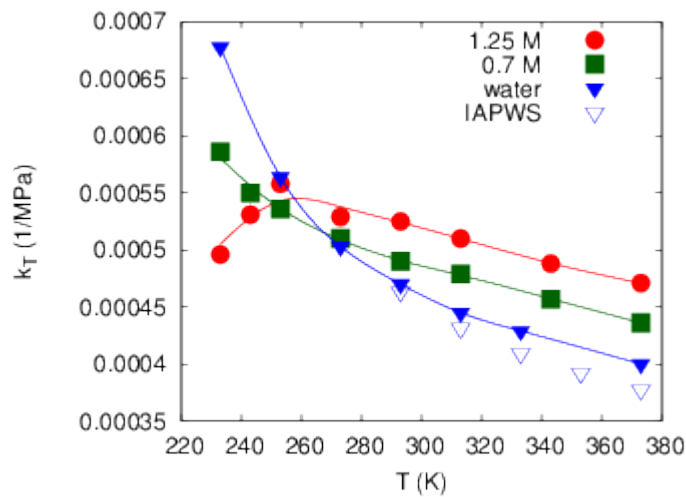
The size of the cluster cluster seems to be independent of the water model



It is impossible that liquid water nucleates homogeneously in Earth for a supercooling $\Delta T < 20 \text{ K}$

Ongoing work on NaCl solutions

- Effect of salt on water anomalies (at extreme conditions)
- Effect of salt on nucleation rate
- Development of a force field for ionic solutions based on TIP4P/2005s



Acknowledgements

We acknowledge the time allocated on the computational facilities Minotauro and Tirant from the Spanish Supercomputing Network (RES) along with the technical support.

UCM team

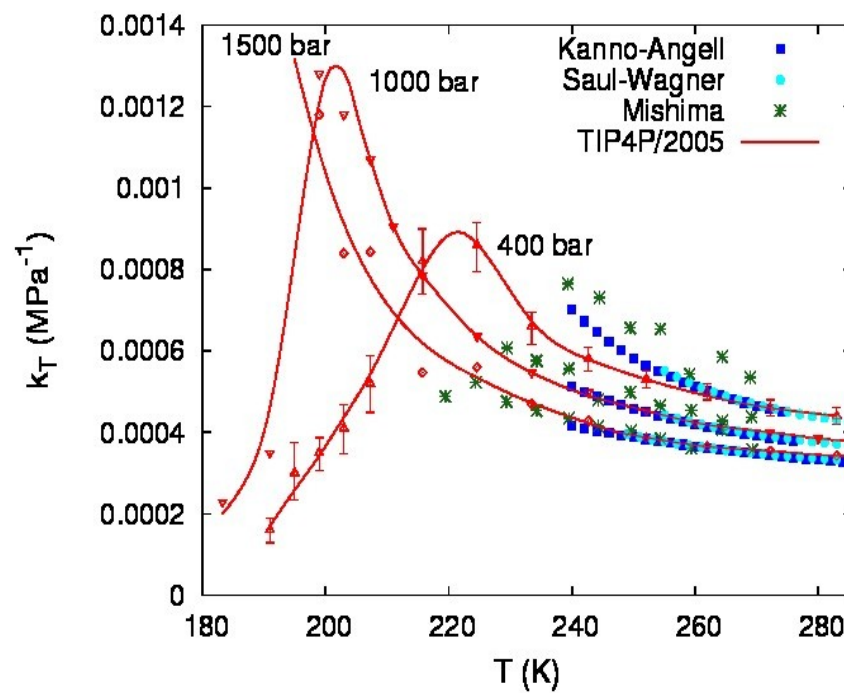
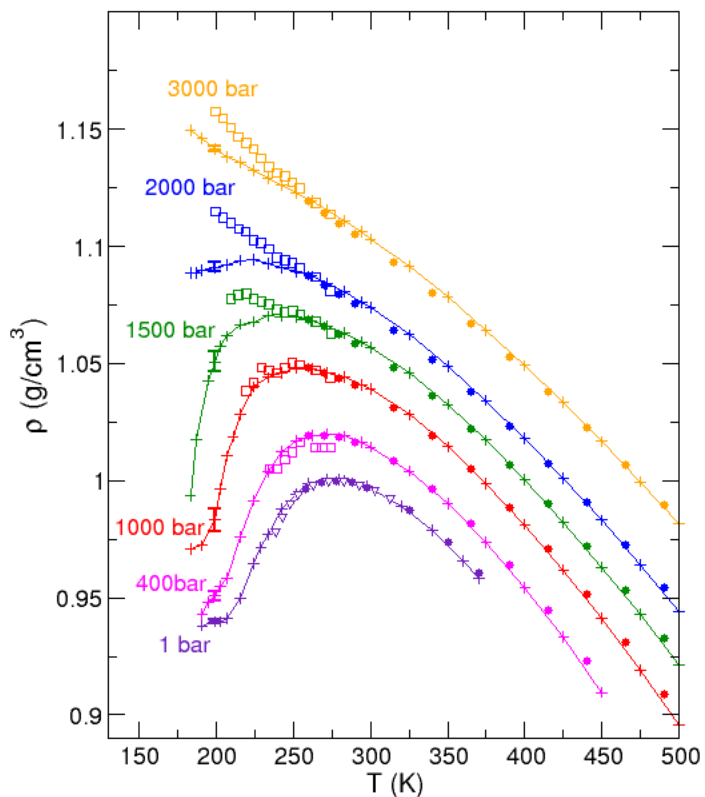
- Carlos Vega
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- Chantal Valeriani
- Miguel Angel González
- Jorge Reñé
- Pablo Rosales
- Miguel Aníbal Portillo
- Guiomar Delgado
- Alberto Zaragoza

Université de Lyon

- Gaël Pallares
- Mouna El Mekki Azouzi
- Frédéric Caupin

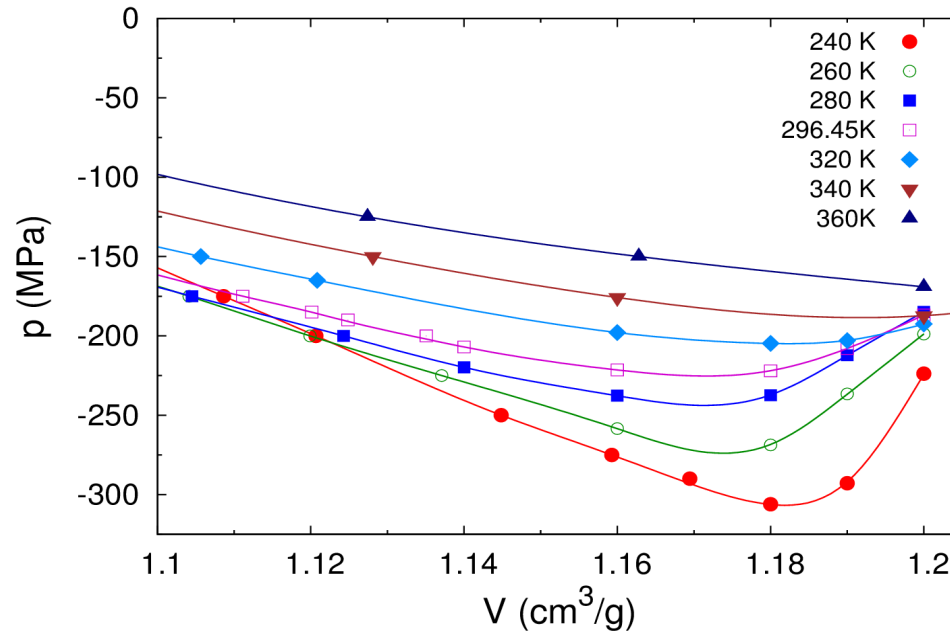


Supercooled TIP4P/2005 water: simulation vs. experiment

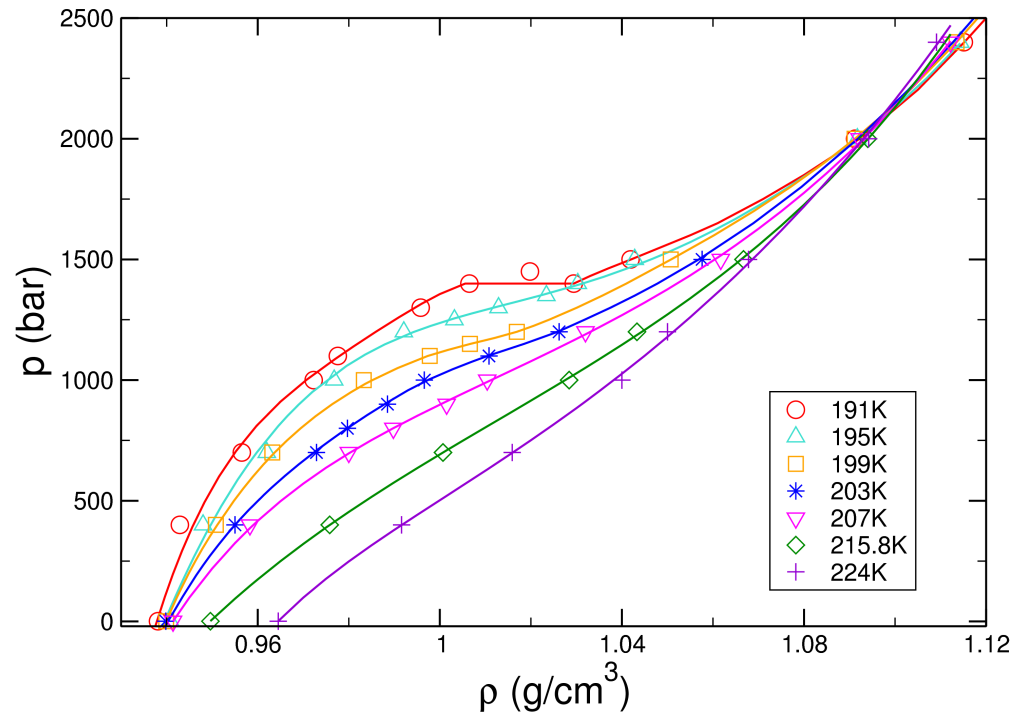


J.L.F. Abascal and C. Vega, JCP 134 186101 (2011)

liquid-vapour spinodal



A liquid-liquid critical point?



J.L.F. Abascal and C. Vega, JCP 133, 234502 (2010)

Cálculo de $\Delta\mu$: Integración

Termodinámica


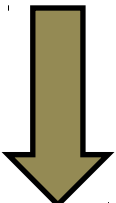
$$\frac{\mu(T_2, p)}{k_B T_2} - \frac{\mu(T_1, p)}{k_B T_1} = - \int_{T_1}^{T_2} \frac{H(T)}{N k_B T^2} dT$$

Hacemos $T_1 = T_{coex}$
y $T_2 = T^*$

$$\Delta\mu \left\{ \begin{array}{l} \frac{\mu(T_2, p)}{k_B T_2} - \frac{\mu(T_1, p)}{k_B T_1} = - \int_{T_1}^{T_2} \frac{H(T)}{N k_B T^2} dT \\ \frac{\mu(T_2, p)}{k_B T_2} - \frac{\mu(T_1, p)}{k_B T_1} = - \int_{T_1}^{T_2} \frac{H(T)}{N k_B T^2} dT \end{array} \right.$$

$$\Delta\mu = [\mu_{fase1}(T_2) - \mu_{fase1}(T_{coex})] - [\mu_{fase2}(T_2) - \mu_{fase2}(T_{coex})] = \mu_{fase1}(T_2) - \mu_{fase2}(T_2)$$

Parámetros de Orden

$$q_{lm}(i) = \frac{1}{N_b(i)} \sum_{j=1}^{N_b(i)} Y_{lm}(\vec{r}_{ij}) \quad \longrightarrow \quad \bar{q}_{lm}(i) = \frac{1}{N_b'(i)} \sum_{j=0}^{N_b'(i)} q_{lm}(j)$$

$$q_l(i) = \sqrt{\frac{4\pi}{2l+1} \sum_{m=-l}^l |q_{lm}(i)|^2} \quad \longrightarrow \quad \bar{q}_l(i) = \sqrt{\frac{4\pi}{2l+1} \sum_{m=-l}^l |\bar{q}_{lm}(i)|^2}$$

$N_b'(i) \equiv \text{Vecinos de la partícula } i + \text{partícula } i$